Classification of missing values in spatial data using spin models

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A problem of current interest is the estimation of spatially distributed processes at locations where measurements are missing. Linear interpolation methods rely on the Gaussian assumption, which is often unrealistic in practice, or normalizing transformations, which are successful only for mild deviations from the Gaussian behavior. We propose to address the problem of missing value estimation on two-dimensional grids by means of spatial classification methods based on spin (Ising, Potts, and clock) models. The "spin" variables provide an interval discretization of the process values, and the spatial correlations are captured in terms of interactions between the spins. The spins at the unmeasured locations are classified by means of the "energy matching" principle: the correlation energy of the entire grid (including prediction sites) is estimated from the sample-based correlations. We investigate the performance of the spin classifiers in terms of computational speed, misclassification rate, class histogram, and spatial correlations reproduction using simulated realizations of spatial random fields, real rainfall data, and a digital test image. We also compare the spin-based methods with standard classifiers such as the *k*-nearest neighbor, the fuzzy *k*-nearest neighbor, and the support vector machine. We find that the spin-based classifiers provide competitive choices.

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I. INTRODUCTION

To date there is a growing demand for various types of spatial data, including remote sensing images, such as thematic maps representing geographical variables, natural resources, land use, environmental indicators, or economic and demographic information. At the same time, new methods are needed for the efficient and reliable processing, analysis, and digital storage of this information. Common issues that arise in the processing phase include the heterogeneity of the data, i.e., the fact that they come from different sources (sensors) operating at different resolutions and often with different biases. In addition, the data coverage is often incomplete due to limited resources (material, human, or technical), equipment limitations (detection level and resolution), meteorological conditions (observations hindered by clouds), and sensor malfunctions.

In the following, we will assume that an observed spatially distributed process in two dimensions is a realization of a spatial random field (SRF) [1]. An SRF, $Z(\vec{r}, \omega)$, where $\vec{r} = (x,y) \in \mathbb{R}^2$ is the location and ω is the state index, represents an ensemble of states (realizations), the spatial correlations of which are determined by a joint probability density function (PDF). The state index will be suppressed in the following to keep the notation concise. In order to use standard tools for the analysis of space-time data, the observed process needs to be estimated at the missing value locations. The spatial estimation can be performed by means of established interpolation and classification techniques [2]. However, considering the ever increasing size of spatial data,

classical methods, e.g., kriging [3] and minimum curvature estimation [4], become impractical due to high computational complexity. Furthermore, the linear spatial interpolation methods assume a jointly Gaussian PDF, which is often not justified by data. In addition, the use of such methods typically requires a considerable degree of subjective input [5]. The nonlinear indicator kriging (IK) method is based on a binary discretization of the data distribution with respect to an arbitrary number of thresholds [6]. IK does not require the Gaussian assumption, but the predicted indicator values are not guaranteed to take 0 or 1 values and may even lie outside the [0, 1] range.

Recent studies investigate applications of statistical physics concepts in various nontraditional fields, such as economy and finance [7–9], materials science [10–12], and biology [13]. Most studies have focused on long-range correlations, while short-range correlations that can be used for spatial or temporal interpolation have received less attention. Nevertheless, within the Gibbs-Markov random field framework, the Potts and Ising models have been used in image analysis [14–18]. The Potts model in the superparamagnetic regime has also been applied to the data clustering problem [19]. In addition, a nonparametric approach based on the Potts model and simulated annealing has been proposed for oil reservoir characterization [20].

Let us consider the Gibbs probability density function $f_Z = e^{-H[Z(\vec{r})]/k_BT}/Z$, where $H[Z(\vec{r})]$ is an interaction functional that involves spatial operators acting on the field $Z(\vec{r})$, k_B is the Boltzmann constant, T is the temperature (we can set $k_BT=1$ without loss of generality), and Z is the partition function. Assume that $H[Z(\vec{r})]$ measures the "spatial disorder energy" of the observed process. Hence, $H[Z(\vec{r})]$ is minimized for a uniform state and increases proportionally to the fluctuations of Z, as prescribed by the interactions embodied in $H[Z(\vec{r})]$. The family of Spartan spatial random fields (SS-RFs) is based on the idea that the spatial correlations can be adequately determined from a general $H[Z(\vec{r})]$ [21] with lo-

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cal interactions. However, the SSRF models are still based on the restrictive Gaussian assumption.

Mild deviations from Gaussian dependence including lognormal dependence can be treated by means of normalizing nonlinear transformations. Alternatively, one can try to incorporate the non-Gaussian terms directly into the statistical model. However, it is not clear what type of interactions in $H[Z(\vec{r})]$ can generate the non-Gaussian behavior encountered in different data sets. Even if such interactions could be defined, calculation of the statistical moments would entail numerical integrations. The purpose of this study is to present nonparametric and computationally efficient methods that do not assume a specific form for the PDF of the observed process. The main idea is to employ interaction models either of the form $H[S^q(Z(\vec{r}))]$, where $S^q[Z(\vec{r})]$ is a binary discretization of the continuous SRF $Z(\vec{r})$ and q is a level index (i.e., Ising model), or of the form $H[S(Z(\vec{r}))]$, where $S[Z(\vec{r})]$ is a multilevel discretization of Z (i.e., Potts and clock models). For the interaction functional we use spin (Ising, Potts, and clock) Hamiltonians. Parameter estimation in such systems is typically based on computationally intensive Monte Carlo (MC) simulations or potentially inaccurate approximation methods. To overcome these problems, we develop a nonparametric Monte Carlo based approach, which benefits from an informative assignment of the initial spin states.

The rest of the paper is organized as follows. In Sec. II, we discuss the problem of spatial classification and prediction and review some established machine learning classification algorithms. In Sec. III we develop our spatial classification approach based on spin models. Section IV investigates the application of the proposed models to synthetic realizations of spatial random fields. Section V focuses on the application of the spin-based classifiers on real data. Finally, in Sec. VI we summarize our results and discuss issues for future research.

II. SPATIAL PREDICTION AND CLASSIFICATION

We consider observations distributed on rectangular grids G of size $N_G = L_x \times L_y$, where L_x and L_y are the horizontal and vertical dimensions of the rectangle (in terms of the unit length), respectively. We denote the set of sampling points by $G_s = \{\vec{r}_i\}$, where $\vec{r}_i = (x_i, y_i) \in \mathbb{R}^2$, $i = 1, \dots, N$ are points scattered on the grid and $N=(1-p)N_G < N_G$. Let z_i be the observation at \vec{r}_i obtained from a joint PDF $f_Z(z_1, \dots, z_{N_C})$. The set $Z_s = \{z_i \in \mathbb{R}\}$ represents the sample of the process. Let G_p $=\{\vec{r}_p\},\ p=1,\ldots,P=pN_G$, be the prediction set so that \tilde{G} $=G_s \cup G_p$. In the following we discretize the continuous distribution of Z using a number of classes (intervals), C_q , q=1,..., N_c . The classes are defined with respect to threshold levels t_k , $k=1,\ldots,N_c+1$. If $Z_{\min}=\min(z_1,\ldots,z_N)$ and Z_{\max} $=\max(z_1,\ldots,z_N)$ denote the minimum and maximum values, respectively, of the data and ϵ is an infinitesimal positive number, then $t_1 = Z_{\min} - \epsilon$ and $t_{N_c+1} = Z_{\max}$. The remaining thresholds are defined by means of $t_q = (q-1)(Z_{\max} - Z_{\min})/N_c + Z_{\min}$ for $q = 2, \dots, N_c$. Each class C_q corresponds to an interval $C_q = (t_q, t_{q+1}]$ for $q = 1, \dots, N_c$. The class indication of the L(z) of ztor field $I_Z(\vec{r})$ is defined so that $I_Z(\vec{r}) = q$ if $z(\vec{r}) \in C_q$ for q $=1,...,N_c$, i.e.,

$$I_{Z}(\vec{r}_{i}) = \sum_{q=1}^{N_{c}} q[\theta(z_{i} - t_{q}) - \theta(z_{i} - t_{q+1})], \quad \forall i = 1, \dots, N,$$
(1)

where $\theta(x)=1$ for x>0 and $\theta(x)=0$ for $x\leq 0$ is the unit step function. Prediction of the field values $Z\{G_p\}$ is then mapped onto a classification problem, i.e., the estimation of $\hat{I}_Z\{G_p\}$. For environmental monitoring and risk management applications useful answers can be obtained in terms of a small number of levels (e.g., eight). For the analysis of gray-scale digital images 256 levels are typically required. As the number of levels tends to infinity, the discretized interval representation tends to the continuum, and spatial classification tends to spatial interpolation.

To test the performance of the spin-based classification methods, we use synthetic Gaussian (normal) and log-normal random field realizations on $L \times L$ grids (where L = 50,100,200), continuous precipitation data sampled on a 50×50 grid, and a digital 512×512 image of 256 gray-scale values. Each data set is randomly and exhaustively partitioned into sampling and prediction sets. We use the points in the prediction set for validation, i.e., to compare the true value of the process with the classifier predictions. Next, we briefly review two nonparametric machine learning classification algorithms, which we use as benchmarks for the performance of the spin models.

A. k-nearest-neighbor models

The k-nearest-neighbor (KNN) model is probably the simplest of all machine learning algorithms [22]. The classified value is assigned to the most populated class among the k nearest neighbors of the classification point. The distance metric used is typically Euclidean, and the neighbors are selected from the points in G_s . The optimal choice of the parameter k is data dependent. Various heuristic techniques can be used to determine the optimal k, e.g., cross validation. However, the method is sensitive to noise or the presence of irrelevant features at the scales of interest. Nevertheless, KNN typically outperforms many other flexible nonlinear methods particularly when the number of explanatory variables is high [23]. It is also easy to implement, and its classification error is asymptotically optimal.

An extension of KNN is the method of fuzzy k-nearest-neighbor (FKNN) classification [24]. In FKNN, the sampling points \vec{r}_j are first assigned a membership to each class C_q , $q=1,\ldots,N_c$, by means of the function $u_q(\vec{r}_j)$. Then, each prediction point \vec{r}_p is also assigned class membership according to the function $u_q(\vec{r}_p) = [\sum_{j=1}^k u_q(\vec{r}_j) || \vec{r}_p - \vec{r}_j ||^{2/(1-m)})$ for $q=1,\ldots,N_c$. The parameter m controls the influence of distant samples. Following [24] we set m=2. The prediction points are classified according to the maxima of the membership functions. The FKNN classifier statistically reduces the effect of noisy samples and produces overall more accurate results than the classical KNN classifier. To eliminate the impact of an arbitrary k, we repeat the classification for $k=1,\ldots,k_{\rm max}$ to determine a $k_{\rm opt}$ that minimizes the misclassification rate. This adaptive ap-

proach guarantees that the lowest misclassification rates achievable by the KNN and the FKNN algorithms (based on the Euclidean-distance metric) are used in the comparisons.

B. Support vector machines

The support vector machine (SVM) classifier is a supervised learning algorithm [25–27] which in several comparative studies outperformed other methods [28–30]. The original SVM algorithm [25] is a linear classifier that segregates data into two classes using maximum-margin hyperplanes. The SVM algorithm has been extended to multiclass and nonlinear problems using the kernel function trick [26] by means of which nonlinear dependence is linearized in a higher-dimensional "feature" Hilbert space. The SVM method is robust and can handle high-dimensional data. However, it is computationally intensive, especially for large N and N_c , because it requires the careful tuning of hyperparameters for each binary classifier followed by the solution of a quadratic problem with N variables.

We solve the $N_c>2$ classification problem by means of N_c binary classifiers operating in one-to-rest fashion, e.g., [31]. We use the software GEOSVM, which implements an adaptation of SVM for spatial data classification [32]. The code is run with radial basis function kernels and involves two tunable hyperparameters: the kernel bandwidth σ_k and the regularization parameter C; the latter controls the trade-off between the machine complexity and the number of linearly nonseparable points. The hyperparameters are tuned to minimize the misclassification rate. Due to the high computational cost of tuning and training the SVM, it is only applied to the rainfall data.

III. SPATIAL CLASSIFICATION BASED ON SPIN MODELS

Below we propose three nonparametric classifiers that use spin model Hamiltonians from statistical physics. In the following, the spins correspond to discretized levels of the continuous variable Z and should not be confused with magnetic moments. The main idea is that the spatial correlations of Z can be captured by the spin interactions. By focusing on the spins it is not necessary to assume a specific form of the joint PDF f_Z . The nonparametric form of the classifiers derives from the fact that the state of the spin systems is constrained by the sample data instead of the interaction couplings J and the temperature T. This is convenient since J and T are unknown a priori, their estimation from the data can be computationally intensive due to the intractability of the partition function, and their uniqueness is not guaranteed for any given sample. To classify the values at the prediction points we use the heuristic principle of energy matching: we calculate the correlation energy of the sample normalized by the number of interacting spin pairs and then determine the spin values at the prediction sites so that the normalized energy for the entire grid matches the respective sample value. Herein we focus on nearest-neighbor correlations, but this constraint can be relaxed.

The idea of correlation energy matching has been applied in the statistical reconstruction of digitized (binary) random media from limited morphological information [33,34]. The classifiers proposed here employ both local (sample values) and global (statistical) information. In particular, this is achieved by performing conditional simulations, in which the sample values are respected locally and the correlation energy globally. This means that while the interactions are translation invariant, the state of the system is not necessarily stationary (statistically homogeneous). The correlation energy matching presupposes that the nearest-neighbor separations in the sample capture the target scale of the prediction set. For example, assume that a sample is drawn from a square grid with step α by removing 50% of the points. The energy matching will be more effective if the points are removed at random than if every second point is removed. In the first case, it is likely that contiguous groups of points will be removed, leaving pairs of neighbors separated by α , while in the second case the minimum separation between the sampling points will be 2α .

The Ising model [35] was introduced to describe the energy states of magnetic materials and later found many applications, e.g., as a model of neuron activity in the brain. It involves binary variables s_i (spins), which can take the values 1 or -1. In the absence of an external filed, the energy of the spin system can be expressed by the Hamiltonian H_I = $-\Sigma_{i,j}J_{ij}s_is_j$. The coupling parameter J_{ij} controls the type $(J_{ij}$ >0 for ferromagnetic and J_{ij} <0 for antiferromagnetic coupling) and strength of the pairwise interactions. The model is usually defined on a regular grid, the interactions are considered uniform, and their range is limited to nearest neighbors. Generalizations to irregular grids and longer-range interactions are also possible. The Ising model has been solved in one dimension and in d=2 without external field. The Potts model is a generalization of the Ising model [36]. Each spin variable is assigned an integer value $s_i \in \{1, ..., N_c\}$, where N_c represents the total number of states. The Hamiltonian of the Potts model is given by $H_P = -\sum_{i,j} J_{ij} \delta(s_i, s_j)$, where δ is the Kronecker symbol. Hence, nonzero contributions to the energy only come from spins in the same state. For $N_c=2$ the Potts model is equivalent to the two-dimensional Ising model. The clock model, also known as the vector Potts model, assumes that the spins take one of N_c possible values, which are distributed uniformly around a circle. The clock Hamiltonian is given by $H_C = -\sum_{i,j} J_{ij} \cos[\frac{2\pi}{N_c}(s_i - s_j)]$. In contrast with the Potts model, in the clock model interactions between spins in different states contribute to the interaction energy. The XY model is obtained from the clock model at the continuum limit $N_c \rightarrow \infty$. The presence of an external field h_i implies an additional term $-\sum_i h_i s_i$ in the Hamiltonian. This term breaks the symmetry of the energy under spin reversals and controls the total magnetization of the spins.

In typical applications of the spin models the parameters J_{ij} and h_i are assumed to be known, and the problem of interest is the estimation of various correlation functions. In the case of spatial data analysis, the parameters are not known *a priori*. Hence, prediction of the spins at unmeasured locations requires us to determine the model parameters from the available spin configuration (sample). The standard procedure for solving such an inverse problem is to infer the parameters by means of the maximum likelihood estimation (MLE) method. Then, the spin values at unsampled locations

can be predicted by maximizing the PDF f_Z (equivalently by minimizing H) under the data constraints. However, the analytical intractability of $\mathcal Z$ hampers the application of MLE. Possible ways of circumventing this problem, such as the maximum pseudolikelihood [37] approach or Markov chain Monte Carlo techniques [38], can be inaccurate or prohibitively slow.

To bypass the problems mentioned above, we propose a nonparametric approach. For lattice data we assume that J_{ij} = J for nearest neighbors and J_{ij} =0 otherwise. In addition, we consider a zero external field. Nevertheless, as shown by the simulation results below, the marginal class distributions of the data are recovered based on the interaction energy of the discretized levels and a judicious choice of the initial states (for each class). The use of different spin models in the study allows us to investigate the impact on the misclassification rate and prediction errors of different classification strategies (sequential vs simultaneous), different couplings (intralevel vs interlevel interactions), as well as the transition to the continuum case (XY model).

A. Classification based on the Ising model

Here we present a spatial classification approach based on the Ising spin model with nearest-neighbor correlations (INNC). The main idea is to use a sequential discretization scheme to estimate the class indicator field $\hat{I}_Z(G_p)$. In this scheme the sample G_s^q and prediction G_p^q grids are progressively updated for increasing class index q. For the lowest class $G_s^1 = G_s$, $G_p^1 = G_p$, where G_s and G_p are the sampling and prediction grids, respectively. Let N_q denotes the number of sites on G_s^q that have fixed class indicator values at level q; at the first level $N_1 = N$ and $N_{q+1} \ge N_q$ for $q = 1, \ldots, N_c - 1$ since the sample at level q + 1 is augmented by the points that are assigned fixed indicator values at level q. Let $S_s^q = \{s_i^q\}, q = 1, \ldots, N_c$; $i_q = 1, \ldots, N_q$ be the set that includes all the sample spin values with respect to level q, and $\widetilde{S}^q = S_s^q \cup S_p^q$ denote all the spin values on \widetilde{G} . The Ising model is used to represent interactions between the spins \widetilde{S}^q .

At each level, the discretization is binary with respect to the upper threshold; i.e., $s_i^q = -1$ if $z_i \le t_{q+1}$ and $s_i^q = 1$ if $z_i > t_{q+1}$ for all $\vec{r}_i \in \tilde{G}$. The classification algorithm sweeps sequentially through the q values. All spin -1 assignments at level q fix the class indicator value for the respective sites; i.e., $\hat{I}_Z(\vec{r}_i) = q$. For q > 1, the sample grid S_s^q is augmented by the points \vec{r}_l for which $s_l^{q-1} = -1$, while the prediction grid is accordingly diminished. It follows that $N_{q>1} \ge N$. The "gaps" in the prediction grid G_p are gradually filled as the algorithm proceeds through consecutive levels. The reduced prediction grid G_p^q at level q contains P^q points so that $P^1 = P$ and $P^q \le P^{q'}$ for q > q'.

The spin assignments at each level are controlled by the cost function, $U(S_p^q|S_s^q)$, which measures the deviation between the correlation energy (per spin pair) of the simulated state, \tilde{C}^q , and the respective sample energy C_s^q . This cost function is given by

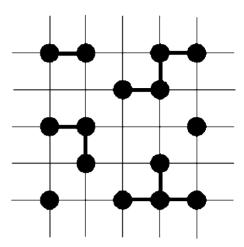


FIG. 1. Schematic depicting the sample sites (solid circles) and the nearest-neighbor pairs (solid circles linked with bold lines) contributing to the sample correlation energy C_s^q .

$$U(S_{p}^{q}|S_{s}^{q}) = \begin{cases} \left[1 - \tilde{C}^{q}(S_{p}^{q}, S_{s}^{q}) / C_{s}^{q}(S_{s}^{q})\right]^{2} & \text{for } C_{s}^{q}(S_{p}^{q}, S_{s}^{q}) \neq 0\\ \left[\tilde{C}^{q}(S_{p}^{q}, S_{s}^{q})\right]^{2} & \text{for } C_{s}^{q}(S_{p}^{q}, S_{s}^{q}) = 0, \end{cases}$$
(2)

where $C_s^q = \langle s_i^q s_j^q \rangle_{G_s^q}$ is the spin two-point correlation of the q-level sample (see Fig. 1) and $\tilde{C}^q = \langle s_i^q s_j^q \rangle_{\tilde{G}}$ is the respective correlation over the entire grid. The q-level classification problem is equivalent to finding the optimal configuration \hat{S}_p^q that minimizes Eq. (2),

$$\hat{S}_p^q = \arg\min_{S_p^q} U(S_p^q | S_s^q). \tag{3}$$

1. Monte Carlo relaxation (greedy scheme)

Determining the optimum spin state is based on a Markov chain MC walk over the ensemble of possible spin configurations at sites that are not assigned a fixed class indicator value. The generation of new states is realized using Metropolis updating at T=0. The zero-temperature condition means that there is no stochastic selection of unfavorable spins; i.e., the spin tested is flipped unconditionally if it lowers the cost function. This so-called "greedy" Monte Carlo algorithm [39] guarantees convergence, which is typically very fast. In contrast, in simulated annealing T is slowly lowered from an initial high-temperature state. This approach is much slower computationally, but the resulting configuration is less sensitive to the initial state. The sensitivity of the greedy algorithm is pronounced in high-dimensional spaces with nonconvex energies since in such cases it is more likely to be trapped in local minima. However, this is not a concern in the current problem. In fact, targeting the global minimum of $U(S_n^q|S_s^q)$ emphasizes exact matching of the sample correlation energy, which is subject to sample-to-sample fluctua-

2. Initial state selection

To provide a fast and unsupervised automatic classification mechanism, the initial configuration of the spins at the prediction sites should require little or no user intervention and minimize the relaxation path to the equilibrium (in state space). Assuming a degree of spatial continuity, which is common in geospatial data sets, we propose to determine the initial states based on the sample values in the immediate neighborhood of the individual prediction points \vec{r}_{p} . More specifically, the initial spin states at the prediction sites are based on majority votes of the sample spins within an adaptable $m \times m$ stencil (where m=2l+1) centered at \vec{r}_n [40]. The stencil size $m \le m_{\text{max}}$ is set adaptively starting with m=3, reflecting the local sampling density and the spin values distribution, to the smallest size that contains a clear majority of spin values. Imposing the arbitrary upper bound $m_{\rm max}$ on the stencil size prevents oversmoothing and increasing the computational load (memory and CPU time). Intuitively, m_{max} should be higher for sparser sampling patterns. If a majority is not achieved for $m \le m_{\text{max}}$, the initial spin value is assigned at random from the sample spin values with the most votes. If a majority cannot be achieved due to a lack of sampling points inside the maximum stencil, the initial value is drawn randomly from the entire range of spin values [41]. We will refer to this procedure for determining the initial state as the adaptable stencil size (ASS) procedure. The indeterminacy of the initial state injects a stochastic element in the classification. Hence, by performing multiple initializations one can assess the classification uncertainty due to the initial state ambiguity.

The parametrization required by the algorithm involves only $m_{\rm max}$ and the definition of the class intervals. The number of classes depends on the study objective: if the goal is to determine areas where a specific level is exceeded, a binary classification is sufficient. For environmental monitoring and

decision making purposes a moderate number of classes (e.g., eight) is often sufficient. A larger number of classes are necessary for the reconstruction of gray-scale images, and even larger numbers of classes are used for spatial interpolation

3. Vectorized algorithm for MC simulation

On commensurate lattices, the grid structure and the short range of the interactions enable vectorization of the algorithm, thus improving the computational efficiency of the Monte Carlo relaxation. Vectorization is achieved by partitioning \tilde{G} in two interpenetrating subgrids \tilde{G}_k , k=1,2, so that the spins on one subgrid interact only with spins on the other. Hence, each subgrid can be updated simultaneously, while one sweep through the entire grid involves just two subgrid updating steps.

Starting from the initial state, Monte Carlo spin updating cycles are generated. Each cycle consists of two steps. Each step focuses on one sublattice, simultaneously generating updates for all the prediction sites. The updates are drawn uniformly from the set of possible values (e.g., ± 1 for the Ising model and $1, \ldots, N_c$ for the Potts and clock models). The updates are accepted or rejected locally for each site, depending on whether they raise or lower the energy of the specific spin's interactions. The algorithm proceeds until the cost function becomes zero within a specified tolerance (termination criterion I) or an updating cycle ends in rejection of all the updates (termination criterion II).

Based on the above, the vectorized MC algorithm for the INNC classification model consists of the following steps:

Algorithm for INNC model

```
(1) Initialization
     (1.1) Define N_c and m_{\text{max}}
     (1.2) Set the indicator field on the entire grid by means of \hat{I}_{z}(\tilde{G})=NaN
(2) Set the simulation level (class index) to q=1
(3) While [loop 1] q \le N_c - 1
     (3.1) Define S_s^q by discretizing Z\{G_s\} with respect to t_{q+1}
     (3.2) Assign spin values S_s^q to the points on G_s^q
     (3.3) Given the data S_s^q, calculate the sample correlation energy C_{I;s}^q
    (3.4) Initialize spin values on G_p^q, based on ASS, i.e., generate \hat{S}_p^{q(0)}
    (3.5) Given S_s^q and \hat{S}_p^{q(0)}, calculate the initial simulated correlation energy \tilde{C}_I^{q(0)}
    (3.6) Initialize the MC index i_{mc}=0 and the rejected states index i_r=0
    (3.7) If \tilde{C}_{I}^{q(0)} < C_{I}^{q}
       (3.7.1) While [loop 1.1] (\tilde{C}_{I}^{q(i_{mc})} < C_{I:s}^{q}) & (i_{r} < 2) update grid
           Initialize the rejected states index i_r=0
           For k=1,2 [subgrid updating]
           (3.7.1.1) Generate a new state \hat{S}_p^{q(i_{mc}+1)} by perturbing \hat{S}_p^{q(i_{mc})} on subgrid \tilde{G}_k
           (3.7.1.2) Identify prediction sites \{i_k\}=K_k\subset \widetilde{G}_k, such that \{s_{i_k}s_j\}^{(i_{mc}+1)}>\{s_{i_k}s_j\}^{(i_{mc})}
           (3.7.1.3) If K_k \neq \emptyset update the sites on K_k
              else keep the "old" state; i_r \rightarrow i_r + 1; endif
           (3.7.1.4) i_{mc} \rightarrow i_{mc} + 1
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end [subgrid updating]
         end [loop 1.1]
      elseif \tilde{C}_{I}^{q(0)} > C_{I}^{q}
         (3.7.2) While [loop 1.2] (\tilde{C}_{I}^{q(i_{mc})} > C_{I:s}^{q}) \& (i_{r} < 2) update grid
             Initialize the rejected states index i_r=0
             For k=1,2 [subgrid updating]
             (3.7.2.1) Generate a new state \hat{S}_p^{q(i_{mc}+1)} by perturbing \hat{S}_p^{q(i_{mc})} on subgrid \tilde{G}_k
             (3.7.2.2) Identify prediction sites \{i_k\}=K_k\subset \widetilde{G}_k, such that \{s_{i_k}s_j\}^{(i_{mc}+1)}<\{s_{i_k}s_j\}^{(i_{mc})}
             (3.7.2.3) If K_k \neq \emptyset update the sites on K_k
                 else keep the "old" state; i_r \rightarrow i_r + 1; endif
             (3.7.2.4) i_{mc} \rightarrow i_{mc} + 1
             end [subgrid updating]
         end [loop 1.2]
      else
        i_{mc} \rightarrow i_{mc} + 1
      (3.8) Assign the -1 spins to the q-th class, i.e.,
         i_{mc} \rightarrow i_{mc} - 1;
         If \hat{S}^{q(i_{mc})}(\{\vec{r}_i\}) = -1, \{\vec{r}_i\} \in \tilde{G}, set \hat{I}_Z(\{\vec{r}_i\}) = q
         where \hat{S}^{q(i_{mc})}(\{\vec{r}_i\} \in G_s) \equiv S_s^q; endif
      (3.9) Increase the class index, q \rightarrow q+1, return to step (3)
end [loop 1]
(4) For q=N_c, \forall \vec{r_i}(i=1,\ldots,N_{\tilde{G}}) such that \hat{I}_Z(\{\vec{r_i}\})=\text{NaN}, set \hat{I}_Z(\{\vec{r_i}\})=N_c.
```

In the above, the symbol NaN denotes a nonnumeric value used to initialize the class assignments. Loop 1 sequentially assigns values to each class. In loop 1.1 the sample energy is approached from below, while in loop 1.2 from above. In both cases the algorithm is vectorized by partitioning the lattice in two subgrids. The termination criterion requires that either the spin correlation energy matches the sample energy (within machine precision) or that one sweep over entire grid not produces any successful updates. In steps 3.7.1.2 and 3.7.2.2 the terms $\{s_{i_l}s_j\}$ imply a summation over all the neighbor spins s_j on the complementary subgrid for each s_{i_k} on the perturbed subgrid. Step 3.8 assigns the -1 spins to the current class and adds the respective sites to the sampling set for the next higher level q. In the end, all the remaining spins with NaN values are assigned to the highest class.

B. Classification based on multivalued spin models

This approach is based on models that involve multivalued spin variables with nearest-neighbor correlations, such

as the Potts, clock, and XY Hamiltonians. The same algorithm structure is used for the nearest-neighbor correlation models based on the Potts (PNNC), clock (CNNC), and XY (XYNNC) models. The PNNC and CNNC models differ only in the form of the correlation energy, while the XYNNC differs from the CNNC model in the number of discretization levels. In contrast with the INNC model, the classification for the multivalued models is performed simultaneously over all levels. Accordingly, we drop the index q which refers to the current level, and the relevant quantities in the algorithm are calculated from the spin values $S = \{s_i\}, i = 1, \dots, N_{\tilde{G}}, \text{ and }$ $s_i \in \{1, \dots, N_c\}$. The normalized correlation energies over the sample and entire grid are calculated by $C_s = \langle \delta(s_i, s_j) \rangle_{G_s}$ and $\begin{array}{lll} \widetilde{C} \!=\! \langle \delta(s_i,s_j) \rangle_{\widetilde{G}} & \text{for the PNNC model and by } C_s \\ \!=\! \langle \cos[\frac{\pi(s_i-s_j)}{2N_c}] \rangle_{G_s} & \text{and } \widetilde{C} \!=\! \langle \cos[\frac{\pi(s_i-s_j)}{2N_c}] \rangle_{\widetilde{G}} & \text{for the CNNC and} \end{array}$ XYNNC models, respectively. Note that the prefactor in the argument of the cosine function is changed from 2π to $\pi/2$ to ensure that each energy level corresponds to a unique value of $|s_i - s_i|$. Based on the above, the MC algorithm for the PNNC models consists of the following steps:

Algorithm for PNNC model

(1) Initialization

- (1.1) Define N_c and m_{max}
- (1.2) Define S_s by discretizing $Z\{G_s\}$ with respect to t_k , $k=1,\ldots,N_c+1$
- (1.3) Assign spin values S_s to the points on G_s

```
(2) Given the data S_s, calculate the sample correlation energy C_{P:s}
(3) Initialize the spin values on G_p, based on ASS, i.e., generate \hat{S}_p^{(0)}
(4) Given S_s and S_p, calculate the initial simulated correlation energy, i.e., \tilde{C}_P^{(0)}
(5) Initialize the \dot{M}C index i_{mc}=0 and the rejected states index i_r=0
(6) If \tilde{C}_{P}^{(0)} < C_{P;s}
     (6.1) While [loop 1] (\tilde{C}_{P}^{(i_{mc})} < C_{P:s}) \& (i_r < 2) update grid
        Initialize the rejected states index i_r=0
        For k=1,2 [subgrid updating]
        (6.1.1) Generate a new state \hat{S}_{p}^{(i_{mc}+1)} by perturbing \hat{S}_{p}^{(i_{mc})} on subgrid \tilde{G}_{k}
        (6.1.2) Identify prediction sites \{i_k\}=K_k\subset \widetilde{G}_k, such that \{\delta(s_{i_k},s_j)\}^{(i_{mc}+1)}>\{\delta(s_{i_k},s_j)\}^{(i_{mc})}
        (6.1.3) If K_k \neq \emptyset update the sites on K_k
            else keep the "old" state; i_r \rightarrow i_r + 1; endif
        (6.1.4) i_{mc} \rightarrow i_{mc} + 1
        end [subgrid updating]
     end [loop 1]
elseif \tilde{C}_{P}^{(0)} > C_{P:s}
     (6.2) While [loop 2] (\tilde{C}_{P}^{(i_{mc})} > C_{P \cdot s}) \& (i_r < 2) update grid
        Initialize the rejected states index i_r=0
        For k=1,2 [subgrid updating]
        (6.2.1) Generate a new state \hat{S}_p^{(i_{mc}+1)} by perturbing \hat{S}_p^{(i_{mc})} on subgrid \tilde{G}_k
        (6.2.2) Identify prediction sites \{i_k\}=K_k\subset \widetilde{G}_k, such that \{\delta(s_i,s_i)\}^{(i_{mc}+1)}<\{\delta(s_i,s_i)\}^{(i_{mc})}
        (6.2.3) If K_k \neq \emptyset update the sites on K_k
            else keep the "old" state; i_r \rightarrow i_r + 1; endif
        (6.2.4) i_{mc} \rightarrow i_{mc} + 1
        end [subgrid updating]
     end [loop 2]
else
     i_{mc} \rightarrow i_{mc} + 1
(7) i_{mc} \rightarrow i_{mc} - 1; assign \hat{I}_{Z}(\{\vec{r}_i\}) = \hat{S}^{(i_{mc})}(\{\vec{r}_i\}), \{\vec{r}_i\} \in \tilde{G}, where \hat{S}^{(i_{mc})}(\{\vec{r}_i\} \in G_s) \equiv S_s.
```

The main difference with the INNC case is the absence of a loop over different classes since the initial spin discretization corresponds to the number of classes.

IV. CLASSIFICATION OF SIMULATED RANDOM FIELD DATA

A. Simulation study design

We study the performance of the classification models on simulated realizations of Gaussian, $Z \sim N(m=50, \sigma=10)$, and log-normal, $\ln Z \sim N(m=4, \sigma=0.5)$, random fields [1]. The spatial correlations are imposed by means of the flexible Whittle-Matérn family of covariance functions,

$$G_{Z}(\|\vec{r}\|) = \sigma^{2} \frac{2^{1-\nu}}{\Gamma(\nu)} (\kappa \|\vec{r}\|)^{\nu} K_{\nu}(\kappa \|\vec{r}\|), \tag{4}$$

where $\|\vec{r}\|$ is the Euclidean two-point distance, σ^2 is the variance, ν is the smoothness parameter, κ is the inverse autocorrelation length, and K_{ν} is the modified Bessel function of index ν . This $G_Z(\|\vec{r}\|)$ leads to random field realizations that are m times differentiable, where m is the largest integer smaller than ν . In addition, higher values of κ imply shorter correlation ranges. We generate samples with different combinations of $\kappa = 0.2, 0.5$ and $\nu = 1.5, 2.5$ on a square $(L_r = L_v)$ =L) grid \tilde{G} , with $N_G = L^2$ nodes, where L = 50, 100, 200. The samples are generated using the spectral method [42,43]. From the complete data set on \tilde{G} , we extract a sample Z_s of size $N=(1-p)N_G$ by randomly removing $P=pN_G$ points for validation. For different degrees of thinning (p=0.33) and 0.66), we generate 100 different configurations of G_n . We use two values of class numbers (N_c =8 and 16), corresponding to different discretizations of the continuum (resolutions). The simulated values on G_p are classified using the spin-based algorithms. In the classification performance evaluation, the indicator field values $I_Z(G_p)$ are compared with the estimates $\hat{I}_Z(G_n)$ that were obtained from the classification algorithms. We define the misclassification rate as

TABLE I. Classification results using the 8- and 16-level INNC, PNNC, and CNNC models for realizations of Gaussian random fields $\sim N(50,10)$ with Whittle-Matérn covariance (κ =0.2, ν =2.5) on a square domain of size L. The tabulated results are averages over 100 realizations. They include the mean value and standard deviation of the misclassification rate $\langle F^* \rangle$ and $\text{STD}_{F^*}(\%)$, the mean number of Monte Carlo sweeps $\langle N_{\text{MC}} \rangle$, the mean optimization CPU time $\langle T_{\text{CPU}} \rangle$ (s), and the mean value of the cost function at termination $\langle U^* \rangle$. The $\langle F^*_{\text{KNN}} \rangle$ and $\langle F^*_{\text{FKNN}} \rangle$ represent the mean of the lowest misclassification rates (%) achieved by KNN and FKNN algorithms, respectively. The number in square brackets indicates the power of 10.

Size			L=	:50					L	=100					L:	=200		
Model	IN	NC	PN	NC	CN	NC	IN	NC	PN	INC	CN	INC	IN	NC	PN	NC	CN	INC
p (%)	33	66	33	66	33	66	33	66	33	66	33	66	33	66	33	66	33	66
									$N_c = 8$									
$\langle F_{\text{KNN}}^* \rangle$	13.6	20.4					10.3	16.1					8.2	12.8				
$\langle F_{ m FKNN}^* \rangle$	12.4	18.1					9.5	14.5					7.6	11.6				
$\langle F^* \rangle$	11.1	16.1	12.9	19.1	12.4	17.5	8.6	12.9	9.6	14.9	9.5	14.0	6.9	10.6	7.4	11.7	7.3	11.1
STD_{F^*}	1.25	1.03	1.34	1.10	1.18	1.02	0.52	0.43	0.56	0.43	0.52	0.45	0.24	0.22	0.26	0.23	0.26	0.21
$\langle N_{ m MC} \rangle$	4.4	6.2	2.7	6.0	3.4	8.5	6.3	6.9	5.5	9.1	5.0	12.5	6.8	8.0	9.5	13.7	9.9	17.4
$\langle T_{\mathrm{CPU}} \rangle$	0.11	0.33	0.03	0.11	0.05	0.13	0.47	1.72	0.13	0.44	0.20	0.61	2.02	7.11	0.61	1.99	1.06	2.95
$\langle U^* \rangle$	5[-4]	2[-3]	3[-5]	1[-4]	8[-9]	4[-8]	3[-4]	1[-3]	5[-6]	2[-5]	2[-9]	1[-8]	1[-4]	6[-4]	1[-6]	4[-6]	3[-10]	2[-9]
									$N_c = 16$	Ó								
$\langle F_{\text{KNN}}^* \rangle$	34.2	40.2					27.7	34.2					20.7	28.0				
$\langle F_{\mathrm{FKNN}}^* \rangle$	31.9	38.9					25.3	32.7					18.9	25.9				
$\langle F^* angle$	21.2	33.8	35.1	42.8	23.5	31.5	17.2	27.3	27.9	35.8	21.2	28.4	13.6	21.5	20.5	28.0	18.2	24.7
STD_{F^*}	1.48	1.60	1.64	1.24	1.52	1.56	0.67	0.75	0.85	0.52	0.73	0.83	0.31	0.36	0.35	0.29	0.38	0.74
$\langle N_{ m MC} \rangle$	8.5	12.9	6.5	17.0	35.1	45.1	13.1	15.2	4.2	14.7	39.6	43.3	13.9	17.1	3.2	2.8	47.5	46.7
$\langle T_{\mathrm{CPU}} \rangle$	0.22	0.62	0.05	0.14	0.18	0.32	0.94	3.30	0.15	0.55	0.84	1.28	3.85	13.92	0.55	1.82	4.02	5.32
$\langle U^* angle$	7[-4]	2[-3]	1[-4]	5[-4]	2[-9]	1[-9]	4[-4]	1[-3]	2[-7]	4[-6]	3[-11]	4[-11]	2[-4]	8[-4]	6[-8]	5[-7]	1[-12]	3[-12]

TABLE II. The same classification statistics are listed as in Table I, obtained from realizations of a Gaussian random field with Whittle-Matérn covariance (κ =0.5, ν =2.5).

Size			L	=50					L:	=100					L:	=200		
Model	IN	NC	PN	NC	CN	NC	IN	NC	PN	NC	CN	NC	IN	NC	PN	NC	CN	INC
p (%)	33	66	33	66	33	66	33	66	33	66	33	66	33	66	33	66	33	66
									$N_c = 8$									
$\langle F_{\mathrm{KNN}}^* \rangle$	25.6	32.3					25.8	32.6					21.8	28.8				
$\langle F_{\rm FKNN}^* \rangle$	23.5	30.3					23.7	30.6					20.2	26.7				
$\langle F^* \rangle$	19.7	27.5	25.6	32.3	22.2	28.7	19.4	27.0	24.9	32.5	22.2	28.8	17.2	23.7	21.2	28.1	19.4	25.5
STD_{F^*}	1.36	1.23	1.63	1.19	1.56	1.33	0.74	0.69	0.85	0.63	0.76	0.71	0.34	0.30	0.43	0.27	0.33	0.29
$\langle N_{ m MC} \rangle$	5.0	6.2	6.6	12.5	7.8	7.9	5.8	7.7	14.6	23.3	6.5	7.9	5.9	8.6	13.0	19.4	6.0	6.6
$\langle T_{\mathrm{CPU}} \rangle$	0.11	0.35	0.04	0.12	0.07	0.14	0.46	1.60	0.18	0.56	0.24	0.53	1.90	6.62	0.74	2.30	0.87	2.07
$\langle U^* angle$	7[-4]	2[-3]	1[-4]	7[-4]	9[-10]	2[-9]	6[-4]	2[-3]	3[-5]	2[-4]	3[-10]	1[-9]	4[-4]	2[-3]	5[-7]	4[-6]	1[-10]	6[-10]
									$N_c = 16$									
$\langle F_{\mathrm{KNN}}^* \rangle$	51.9	55.4					53.9	56.5					48.1	51.4				
$\langle F_{\rm FKNN}^* \rangle$	49.7	54.7					51.5	56.1					45.4	50.9				
$\langle F^* \rangle$	39.0	54.0	52.7	58.5	37.2	48.7	38.7	53.9	52.9	59.3	36.2	48.1	33.7	48.1	46.9	54.1	32.7	44.1
STD_{F^*}	1.77	1.55	1.55	1.00	2.30	1.69	0.87	0.76	0.77	0.54	1.34	0.96	0.47	0.42	0.41	0.24	0.55	0.44
$\langle N_{ m MC} \rangle$	10.3	12.7	19.1	27.3	23.1	19.9	11.6	14.3	39.7	50.1	23.3	20.6	11.9	16.8	59.6	70.4	24.4	21.5
$\langle T_{ m CPU} \rangle$	0.21	0.65	0.07	0.17	0.14	0.21	0.91	2.99	0.34	0.88	0.57	0.85	3.78	12.44	1.90	4.69	2.41	3.57
$\langle U^* angle$	8[-4]	2[-3]	2[-3]	2[-2]	9[-11]	3[-10]	7[-4]	2[-3]	2[-3]	2[-2]	2[-11]	1[-10]	5[-4]	2[-3]	3[-4]	7[-3]	1[-11]	7[-11]

TABLE III. The same classification statistics are listed as in Table I, obtained from realizations of a Gaussian random field with Whittle-Matérn covariance (κ =0.5, ν =1.5).

Size			L=	=50					L=	:100					L=	:200		
Model	IN	NC	PN	NC	CNI	NC	IN	NC	PN	NC	CN	NC	IN	NC	PN	NC	CN	NC
p (%)	33	66	33	66	33	66	33	66	33	66	33	66	33	66	33	66	33	66
									N_c =8									
$\langle F_{\mathrm{KNN}}^* \rangle$	38.9	44.8					36.7	42.4					30.7	37.1				
$\langle F_{\mathrm{FKNN}}^* \rangle$	36.7	42.9					34.6	40.6					28.9	34.9				
$\langle F^* \rangle$	31.7	39.5	38.6	45.0	35.9	42.8	29.5	37.2	36.0	42.4	33.9	40.2	25.6	32.0	29.7	36.1	29.9	35.4
STD_{F^*}	1.66	1.31	1.59	1.14	1.71	1.10	0.79	0.68	0.69	0.57	0.78	0.75	0.37	0.32	0.36	0.28	0.39	0.30
$\langle N_{ m MC} \rangle$	3.9	5.5	12.2	16.9	3.4	3.2	4.1	6.6	21.9	27.5	3.2	3.2	4.2	7.3	29.6	37.2	0.76	1.3
$\langle T_{\mathrm{CPU}} \rangle$	0.12	0.36	0.05	0.14	0.05	0.11	0.48	1.62	0.22	0.64	0.18	0.45	1.98	6.76	1.05	2.86	0.53	1.69
$\langle U^* angle$	9[-4]	3[-3]	3[-3]	1[-2]	4[-9]	2[-8]	7[-4]	3[-3]	1[-3]	7[-3]	2[-9]	1[-8]	4[-4]	2[-3]	8[-4]	4[-3]	1[-9]	8[-9]
									$N_c = 16$									
$\langle F_{\mathrm{KNN}}^* \rangle$	64.8	67.3					63.5	65.7					58.5	61.0				
$\langle F_{\rm FKNN}^* \rangle$	63.2	66.6					61.4	65.1					56.3	60.5				
$\langle F^* \rangle$	54.8	65.8	65.1	69.4	56.7	64.5	53.3	64.8	62.2	67.3	54.8	62.9	47.9	58.4	57.3	62.7	50.8	58.6
STD_{F^*}	1.71	1.34	1.46	1.07	1.77	1.24	0.89	0.76	0.78	0.54	0.89	0.67	0.44	0.34	0.42	0.28	0.50	0.37
$\langle N_{ m MC} \rangle$	8.6	11.0	21.0	29.8	9.2	8.2	8.7	12.4	40.5	51.9	9.2	7.9	8.9	13.8	61.3	73.2	8.9	7.8
$\langle T_{\mathrm{CPU}} \rangle$	0.23	0.67	0.07	0.18	0.09	0.15	0.94	3.09	0.36	0.91	0.32	0.61	3.96	12.78	1.84	4.76	1.28	2.49
$\langle U^* angle$	8[-4]	2[-3]	2[-2]	1[-1]	9[-10]	4[-9]	7[-4]	2[-3]	1[-2]	9[-2]	6[-10]	3[-9]	5[-4]	2[-3]	8[-3]	6[-2]	3[-10]	2[-9]

TABLE IV. The same classification statistics are listed as in Table I, obtained from realizations of a log-normal random field with Whittle-Matérn covariance (κ =0.5, ν =2.5).

Size			L=	=50					L:	=100					L=	=200		
Model	IN	NC	PN	NC	CN	INC	IN	NC	PN	NC	CN	NC	IN	NC	PN	NC	CN	INC
p (%)	33	66	33	66	33	66	33	66	33	66	33	66	33	66	33	66	33	66
									$N_c = 8$									
$\langle F_{ m KNN}^* \rangle$	18.7	24.3					21.5	27.5					16.5	22.3				
$\langle F_{ ext{FKNN}}^* \rangle$	17.3	22.3					19.9	25.4					15.3	20.3				
$\langle F^* \rangle$	15.2	20.9	18.4	23.8	16.3	21.6	16.8	23.1	20.6	26.6	18.2	24.0	13.7	18.8	15.6	20.9	14.9	19.8
STD_{F^*}	1.29	1.17	1.31	0.95	1.15	1.13	0.70	0.63	0.76	0.52	0.67	0.61	0.30	0.30	0.35	0.31	0.31	0.31
$\langle N_{ m MC} \rangle$	3.7	4.6	6.0	10.0	7.0	8.4	4.5	5.7	15.0	20.3	6.9	7.1	4.6	6.5	19.5	20.6	5.1	5.3
$\langle T_{\mathrm{CPU}} \rangle$	0.08	0.19	0.04	0.12	0.07	0.14	0.32	0.87	0.18	0.57	0.24	0.53	1.26	3.27	0.86	2.34	0.83	2.28
$\langle U^* \rangle$	6[-4]	1[-3]	1[-4]	7[-4]	6[-9]	7[-9]	6[-4]	2[-3]	4[-5]	2[-4]	2[-10]	9[-10]	3[-4]	1[-3]	4[-6]	8[-6]	8[-11]	6[-10]
									$N_c = 16$									
$\langle F_{ m KNN}^* \rangle$	37.1	41.7					42.7	47.1					35.1	40.2				
$\langle F_{\rm FKNN}^* \rangle$	34.9	40.0					40.4	45.8					33.1	38.7				
$\langle F^* \rangle$	28.8	37.9	37.0	42.4	29.7	37.1	32.0	42.2	41.7	48.5	31.5	41.4	26.1	35.1	34.0	40.8	27.2	35.6
STD_{F^*}	1.85	1.43	1.62	1.07	1.73	1.41	0.96	0.79	0.75	0.56	1.15	0.74	0.35	0.40	0.37	0.30	0.46	0.42
$\langle N_{ m MC} \rangle$	8.9	10.4	13.8	21.4	18.8	18.8	10.3	12.5	35.7	45.3	19.9	17.5	10.8	14.1	52.1	63.9	19.7	17.2
$\langle T_{ m CPU} \rangle$	0.16	0.33	0.06	0.16	0.12	0.20	0.61	1.53	0.31	0.84	0.54	0.84	2.35	5.7	1.63	4.31	2.14	3.44
$\langle U^* angle$	7[-4]	2[-3]	8[-4]	6[-3]	2[-9]	2[-10]	7[-4]	2[-3]	7[-4]	6[-3]	3[-11]	2[-10]	5[-4]	2[-3]	1[-4]	2[-3]	9[-12]	7[-11]

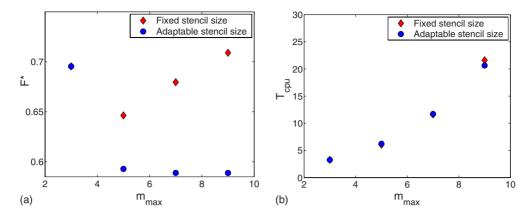


FIG. 2. (Color online) Dependence of misclassification rate, F^* , on the maximum stencil size m_{max} used in the initialization with both adaptable and fixed stencil size (a). Dependence of the CPU time on m_{max} (b). The 16-class INNC model is applied to a sample generated from a realization of a Gaussian random field with Whittle-Matérn covariance (κ =0.5, ν =1.5) on a square grid (L=200) by removing 66% of the points.

$$F = \frac{1}{P} \sum_{p=1}^{P} \{1 - \delta[I_Z(\vec{r}_p), \hat{I}_Z(\vec{r}_p)]\},$$
 (5)

where $I_Z(\vec{r}_p)$ is the true value at the validation points, $\hat{I}_Z(\vec{r}_p)$ is the classification estimate, and $\delta(I,I')=1$ if I=I', $\delta(I,I')=0$ if $I\neq I'$. We also measure the optimization CPU time, T_{CPU} , the number of MC sweeps required for reaching equilibrium, and the residual values of the cost function at termination, U^* .

To evaluate the performance of the spin models for large N_c (i.e., approaching the continuum limit), we calculate the following "sample" prediction errors: average absolute $\operatorname{error} = \frac{1}{p} \sum_{p=1}^P |Z(\vec{r}_p) - \hat{Z}(\vec{r}_p)|$, average relative $\operatorname{error} = \frac{1}{p} \sum_{p=1}^P \frac{Z(\vec{r}_p) - \hat{Z}(\vec{r}_p)}{Z(\vec{r}_p)}$, average absolute relative $\operatorname{error} = \frac{1}{p} \sum_{p=1}^P |\frac{Z(\vec{r}_p) - \hat{Z}(\vec{r}_p)}{Z(\vec{r}_p)}|$, root average square $\operatorname{error} = \sqrt{\sum_{p=1}^P \frac{1}{p}} |Z(\vec{r}_p) - \hat{Z}(\vec{r}_p)|^2$, and the linear sample correlation coefficient R. In the above definitions, P is the number of validation points, $Z(\vec{r}_p)$ is the true value at \vec{r}_p , and $\hat{Z}(\vec{r}_p)$ is the estimate of the process based on

$$\hat{Z}(\vec{r}_p) = t_{\hat{I}_Z(\vec{r}_p)} + \frac{t_{\hat{I}_Z(\vec{r}_p)+1} - t_{\hat{I}_Z(\vec{r}_p)}}{2}.$$
 (6)

focus on the local behavior of the classifiers respective errors, in calculated over a number is of (e.g., M=100 realizations) at each point. Namely, at define \vec{r}_p , $\begin{array}{lll} & \operatorname{error} = \langle \left| Z(\vec{r}_p) - \hat{Z}(\vec{r}_p) \right| \rangle, & \text{the mean relative error} \\ = \langle \frac{Z(\vec{r}_p) - \hat{Z}(\vec{r}_p)}{Z(\vec{r}_p)} \rangle, & \text{the mean absolute relative error} = \langle \left| \frac{Z(\vec{r}_p) - \hat{Z}(\vec{r}_p)}{Z(\vec{r}_p)} \right| \rangle, \end{array}$ and the root mean square error= $\sqrt{\langle |Z(\vec{r}_p) - \hat{Z}(\vec{r}_p)|^2 \rangle}$, where $\hat{Z}(\vec{r}_p)$ are the predictions at point \vec{r}_p and $\langle \cdot \rangle$ denotes the average over M realizations. The computations are performed in the Matlab® programming environment on a desktop computer with 1.93 GB of random access memory and an Intel®CoreTM2 CPU 6320 processor with an 1.86 GHz clock.

B. Misclassification rate

We investigate the effects of grid size N_G , data sparseness p, and class number N_c on the classification performance. The results obtained by the INNC, PNNC, and CNNC mod-

TABLE V. Classification results using the 8- and 16-level INNC, PNNC, and CNNC models for realizations of a Gaussian random field $\sim N(50,10)$ with Whittle-Matérn covariance (κ =0.5, ν =2.5) and L=200. The simulations start from random initial states. The same statistics are listed as in Table I. The number in square brackets indicates the power of 10.

			N_{c}	c=8					N_{c}	=16		
Model	IN	NC	PN	NC	CN	NC	IN	NC	PN	NC	CN	NC
p (%)	33	66	33	66	33	66	33	66	33	66	33	66
$\langle F^* angle$	34.3	91.3	25.6	42.5	22.0	36.3	67.1	98.8	56.5	69.8	32.7	49.4
STD_{F^*}	0.50	0.20	0.34	0.43	0.37	0.47	18.3	0.07	0.40	0.43	0.59	0.64
$\langle N_{ m MC} \rangle$	7.0	3.0	27.3	74.9	31.7	68.9	13.1	5.2	23.3	32.3	42.4	49.8
$\langle T_{\mathrm{CPU}} \rangle$	1.15	1.24	0.63	2.54	2.12	4.92	2.28	2.0	0.53	1.16	2.82	3.73
$\langle U^* \rangle$	1[-1]	1[-1]	2[-7]	8[-4]	2[-11]	3[-8]	9[-2]	6[-2]	4[-6]	3[-6]	9[-12]	4[-11]

TABLE VI. Degeneracy and residual values of the cost function corresponding to 100 random initial states. The sample is a realization of a Gaussian random field $\sim N(50,10)$ with Whittle-Matérn covariance (κ =0.5, ν =2.5) on a grid of size L=50. The number in square brackets indicates the power of 10.

			N_{c}	=8					N	c=16		
Model	IN	NC	PN	NC	CN	NC	IN	NC	PN	NC	CN	NC
p (%)	33	66	33	66	33	66	33	66	33	66	33	66
Degeneracy	6	8	77	10	64	7	8	4	61	53	6	0
$\langle U^* angle$	8[-2]	1[-1]	7[-7]	4[-4]	1[-10]	7[-9]	8[-2]	6[-2]	4[-6]	5[-6]	5[-11]	1[-10]

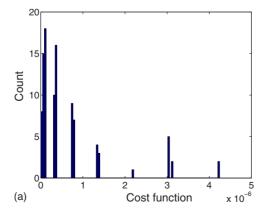
els are compared with the *best* results obtained by the established KNN and FKNN methods (see Sec. II A). Tables I–III summarize the results obtained for Gaussian random fields with Whittle-Matérn parameters $(\kappa, \nu) = (0.2, 2.5)$, (0.5, 2.5), and (0.5, 1.5), respectively. Table IV lists the results obtained from log-normal random fields with Whittle-Matérn parameters $(\kappa, \nu) = (0.5, 2.5)$.

As expected, the misclassification rate overall decreases with increasing N_G and increases proportionally with p. Comparing the performance of different models at fixed N_G , p, N_c , and distributional assumptions, the INNC model performs uniformly better for N_c =8. For N_c =16 and (κ, ν) =(0.2,2.5), the INNC model gives in general the lowest $\langle F^* \rangle$, except one case in which the CNNC model performs best. Generally, the CNNC model is expected to perform better when the cross-class correlations make finite contributions to the correlation energy. This is likely as the number of classes or the spatial variation in the data increase. The case of $N_c = 16$ and $(\kappa, \nu) = (0.5, 2.5)$ provides an example combining higher N_c and shorter-range variations than N_c =8 and κ =0.2. This example shows a clear advantage of the CNNC model (especially at p=0.66) over the others. However, since the interactions are restricted to nearest neighbors, crossclass correlations in CNNC can increase the sensitivity to noisy or "rough" data. This is evidenced in the N_c =16 case for $(\kappa, \nu) = (0.5, 1.5)$. The random field realizations for ν =1.5 are only once differentiable, in contrast with the ν =2.5 case, where the realizations are twice differentiable. For ν =1.5 the classification performance of the CNNC model approaches that of the INNC model. The highest misclassification rate among the spin models investigated is displayed consistently by the PNNC model. This could be attributed to the fact that it incorporates less information than its counterparts. Namely, the INNC model sequentially uses the lower level classification results in the higher level estimates. On the other hand, the CNNC model differentiates between neighbors of various classes. In contrast, the PNNC model can only distinguish if the neighbors belong to the same or different classes. As we show below, the Potts model leads to degenerate states, which are undesirable for classification purposes.

C. Impact of initial state selection

Figure 2 illustrates the behavior of the misclassification rates, F^* , and the CPU times versus $m_{\rm max}$, considering fixed stencil size (FSS) and ASS approaches. Relatively small $m_{\rm max}$ =5,7 for FSS and ASS, respectively, lead to the lowest F^* values. The computational time increases quadratically with $m_{\rm max}$. There is no significant difference in computational time between the two methods. However, using ASS leads to a significantly lower misclassification rate.

For the results obtained in Tables I–IV, the initial state was based on majority vote with ASS. This accelerates the relaxation process and prevents trapping in a "poor" local minimum. To test the sensitivity of the spin models on different initial states, we repeat the simulations using randomly assigned initial states for Gaussian data with $(\kappa, \nu) = (0.5, 2.5)$ on a L=200 grid. The results for different models and two values of N_c are given in Table V. The greatest



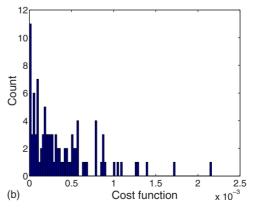


FIG. 3. (Color online) Histograms of the residual values of the cost function obtained from the PNNC model for one sampling realization and 100 different random initial states, with N_c =8 and different values of p=0.33 (a) and p=0.66 (b). Each bar may include states with similar but not identical values of the cost function due to finite bar width. Note different scales of the abscissas.

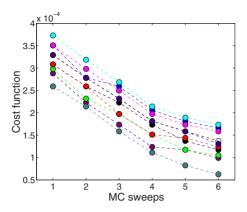


FIG. 4. (Color online) Cost function evolution versus the number of full (over the entire grid) MC sweeps. Ten initial states are generated by the majority rule with adaptable stencil size, leading to ten U^* evolution curves. The curves are based on the 8-class PNNC model simulations performed on a single realization of a Gaussian random field (κ =0.5, ν =2.5) on a grid of size L=50 from which 33% of the points are randomly removed.

impact on the classification performance is observed for the INNC model. On the other hand, the CNNC model is the most robust with respect to changes in the initial state, especially for lower p and higher N_c . For instance, at p = 0.33 and $N_c = 16$, there is practically no difference between the misclassification rates obtained using initial states based on ran-

dom versus ASS majority rule selection of the spins. Using a random initial state increases by 73% the number of MC sweeps needed to achieve equilibrium. However, the CPU time is increased only by 17% because the random assignment strategy leads to faster determination of the initial state.

D. Spatial degeneracy

The higher sensitivity of the INNC model to the initial state is due to the sequential spin estimation, which propagates the misclassification from lower to higher levels, thus resulting in a higher overall misclassification rate. The misclassification at lower levels can occur due to the presence of degenerate states which correspond to different spatial configurations even though their energies are numerically identical. Generally, the degeneracy increases with p, as a result of relaxing the spatial constraints. As the size of the prediction set is reduced at higher levels, due to the inclusion of the classified lower levels in the sampling set, the degeneracy is accordingly diminished. However, the lower level misclassifications propagate to higher levels. A high level of degeneracy is also responsible for the relatively poor classification performance of the PNNC model; since the latter does not include cross-class correlations, the energy contributions of all the nearest-neighbor pairs that do not belong in the same class are zero. We believe that the robustness of the CNNC model and its competitive classification performance is par-

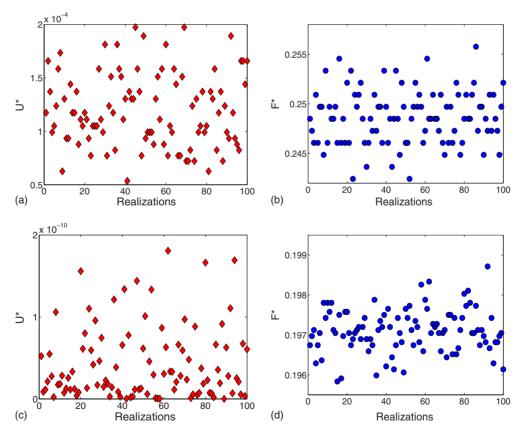


FIG. 5. (Color online) Residual values of the cost function at termination, U^* , and the misclassification rate, F^* , for the 8-class PNNC model [(a) and (b)] and for the 8-class CNNC model [(c) and (d)]. The sample is obtained from one realization of a Gaussian random field with Whittle-Matérn parameters (κ =0.5, ν =2.5) on grids of size L=50 (PNNC) and L=200 (CNNC), respectively, by randomly removing p=0.33 of the points. The U^* , F^* values are obtained from 100 initial states generated by the majority rule with adaptable stencil size.

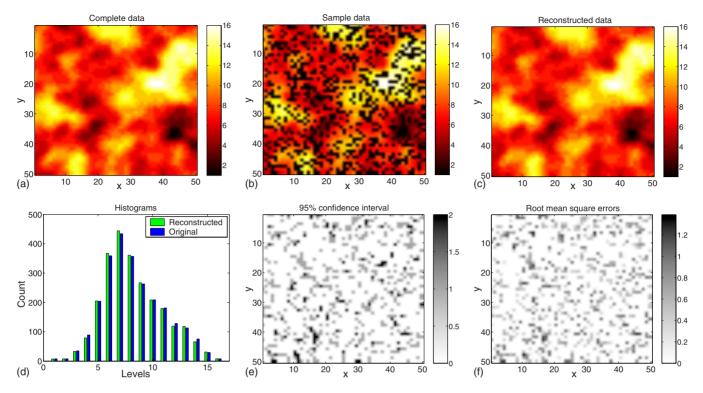


FIG. 6. (Color online) Classification results obtained from the 16-class CNNC model. The sample is obtained from one realization of a Gaussian random field with Whittle-Matérn parameters (κ =0.5, ν =2.5) on a grid of size L=50 by randomly removing p=0.33 of the points. The 100 reconstructions are generated starting from 100 initial states obtained by the majority rule with adaptable stencil size. Plots (a)–(c) show class maps for the complete realization, the sample with the missing data, and the reconstructed image, respectively. The latter is based on the medians of the class values obtained from the 100 reconstructions. Plot (d) compares the class histograms of the original and reconstructed data, plot (e) shows the width of the 95% confidence intervals for the class predictions, and plot (f) represents the class root-mean-square error at the prediction points.

tially due to the reduction in degeneracy achieved by differentiating between the energy contributions of neighbors belonging to different classes.

To investigate the degeneracy of the final class configurations obtained by the different spin models, we use the same sampling configuration on a grid of size L=50. We then run the MC simulation for 100 different (randomly assigned) sets of initial values at the prediction points. We measure the degeneracy as the number of final configurations (states) with the same residual cost function value as at least one other state from which they differ in the misclassification rate. The results are summarized in Table VI. Note that for the INNC model, the degenerate states only appear at the first (q=1) level. The highest degeneracy is observed in the PNNC model. The histograms of the residual cost function value obtained for the PNNC model at p=0.33 and p=0.66are shown in Fig. 3. Intuitively, one would expect the degeneracy to increase with p. However, the opposite tendency is observed in the simulations. The overall shift of $\langle U^* \rangle$ to higher values and the shape of the histogram [Fig. 3(b)] suggest that the reduction in degeneracy is due to a scattering of energy levels. This can be viewed as a result of the cost function developing multiple local minima (multimodality). A reduction in the degeneracy is also observed with increasing L: for example, the configurations produced by the CNNC model for L=200 do not exhibit spatial degeneracy for any of the N_c and p combinations tested.

E. Classification uncertainty

Multimodality and degeneracy introduce an uncertainty in the classification outcome, starting from different initial states. The cost function multimodality persists even in cases where the degeneracy vanishes. An additional source of uncertainty is termination criterion II: if the relaxation process stops before the (local) optimum is reached, the final configuration depends on the initial state, thus mimicking multimodality of the cost function. Termination criterion II is arbitrary and aims at computational efficiency. Hence, it does not guarantee that the resulting configuration is in a perfect (quasi)equilibrium state. Uncertainty is primarily generated by the ASS random initial assignment of those prediction points for which majority rule is not established. Figure 4 demonstrates the cost function evolution during the relaxation process. We use the 8-class PNNC model classification on a single realization of a Gaussian random field (κ =0.5, ν =2.5) on a grid of size L=50 (high spatial degeneracy) thinned by 33%. The ten curves correspond to initial states generated by the majority rule with ASS. Different initial states follow different relaxation paths, leading to different local minima (based on the current value of termination criterion II). Figures 5(a) and 5(b) are the residual values of the cost function at termination and the misclassification rates, respectively, for the 8-class PNNC model, obtained from 100 initial states generated by the majority rule with ASS. Figures 5(c) and 5(d) show the same quantities as Figs.

TABLE VII. Validation statistics of "discrete-level interpolation" obtained by the CNNC model for samples of Gaussian and log-normal random fields on grids of size L=100 with Whittle-Matérn parameters κ =0.5, ν =2.5. MAAE: mean average absolute error. MARE: mean average relative error. MAARE: mean average absolute relative error. MRASE: mean root average square error. MR: mean correlation coefficient. First, averages are evaluated over the prediction points for each realization, and then means are calculated over an ensemble of 100 realizations. The number in square brackets indicates the power of 10.

Data			Normal	mal					Log n	Log normal		
No. of classes	N_c =	$N_c = 32$	N_c =	^r _c =64	$N_c = 128$	128	$N_c = 32$	=32	N_c =	$N_c = 64$	$N_c =$	$N_c = 128$
(%) d	33	99	33	99	33	99	33	99	33	99	33	99
MAAE	1.32	1.96	1.24	1.93	1.21	1.93	11.88	16.97	10.97	16.51	10.73	16.51
MARE (%)	-0.28	-0.49	-0.29	-0.53	-0.30	-0.54	-0.39	-0.12	-0.58	-0.39	-0.81	-0.84
MAARE (%)	2.78	4.14	2.60	4.07	2.56	4.08	7.78	10.68	7.09	10.33	6.95	10.42
MRASE	1.71	2.56	1.61	2.53	1.59	2.54	16.12	23.85	15.10	23.45	14.86	23.41
MR (%)	98.46	96.48	99.86	69.96	98.71	96.61	76.76	95.60	98.25	95.79	98.31	95.79
$\langle N_{ m MC} angle$	33.6	28.3	37.1	31.1	38.7	32.7	28.8	24.7	33.0	27.9	34.8	29.5
$\langle T_{\mathrm{CPU}} \rangle$ (s)	0.78	1.13	1.03	1.50	1.24	2.04	0.76	1.14	0.97	1.35	1.09	1.87
$\langle U^* angle$	8[-12]	7[-11]	7[-12]	4[-11]	6[-12]	5[-11]	1[-11]	7[-11]	7[-12]	6[-11]	7[-12]	5[-11]

5(a) and 5(b) but for L=200 using the 8-class CNNC model (for which no degeneracy was observed). In both cases, the resulting values vary, with the variations being more pronounced in the degenerate case.

In light of the above, the spin-based classification methods permit the estimation of the classification uncertainty by sampling over different initial configurations. This allows us to determine points in space with increased chance of misclassification. Therefore, multiple Monte Carlo runs starting from different initial states permit statistical estimates. This repetitive application can also improve estimation results compared to a single run. For example, we consider one sample realization of a Gaussian random field with κ =0.5, ν =2.5 on a grid of size L=50 reduced by p=0.33. A single simulation run using the 16-class CNNC model gave a misclassification rate of $F^*=36.6\%$. Repeating the simulation 100 times, the median misclassification rate drops to F^* =30.1%. The total CPU time increases linearly with the number of runs, leading to $T_{\rm CPU}$ =14.8 s. We note that comparable results can also be obtained with significantly fewer simulation runs, e.g., 10. Figure 6(a) shows the complete realization, Fig. 6(b) shows the sample with the missing data points, and Fig. 6(c) shows the reconstructed image based on the medians of the estimates from the 100 runs. There is good visual agreement between the original and the reconstructed images. The class histogram of the reconstruction also matches satisfactorily that of the data, both shown in Fig. 6(d). Note that both histograms are asymmetric even though the random field values are normally distributed. This is due to the relatively long correlation length that results in relatively more areas of higher than lower values as shown in Fig. 6(a). Finally, Fig. 6(e) displays maps representing the width of the 95% class confidence interval and Fig. 6(f) displays the class root-mean-square error at each prediction point.

F. Computational efficiency

Thanks to the vectorization of the algorithm and judicious choice of the initial state, all the models perform the classification task very efficiently. The mean CPU time for one realization ranged from 0.03 s for the 8-class PNNC model, with L=50 and p=0.33, up to almost 14 s for the 16-class INNC model, with L=200 and p=0.66. The INNC displayed the highest CPU times in spite of the fact that the equilibration on the respective levels was extremely fast due to the binary nature of the data. The major part of the total CPU time was spent on the initial state assignments, which is repeated at every level unlike the other models. This results in an approximately linear increase in the INNC CPU time with N_c . The best energy matching, marked by the lowest values of $\langle U^* \rangle$, was observed in the CNNC model, which can closely approximate the sample energy due to the flexibility of the multilevel interactions. In general, the poorest matching is shown by the PNNC model at high N_c and p, indicating a high multimodality of the cost function and a failure to reach a "good" local optimum. In comparison with the reference models, the INNC and CNNC models gave systematically better classification performance than the best results

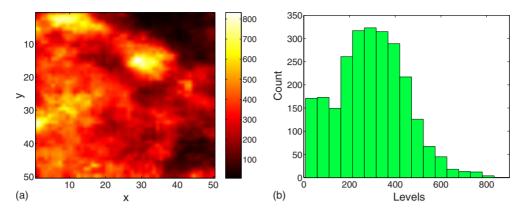


FIG. 7. (Color online) Map and histogram of precipitation values for the complete rainfall data.

obtained by the KNN and the FKNN algorithms. The PNNC model showed systematically worse classification performance than the FKNN, but it gave better results than the regular KNN for relatively small N_c and p.

G. Approaching the continuum class limit

If N_c is high, the classification problem becomes a regression problem. It is then relevant to monitor quantities such as prediction errors and the correlation coefficient between the true (not discretized) and the predicted values back transformed to the original data scale. Table VII focuses on the performance of the CNNC model with gradual increase in number of classes up to N_c =128 (extension to XYNNC model) for the selected size and types of data. The errors (except for MARE) show a decreasing tendency while the mean correlation coefficient MR increases. These tendencies seem to persist up to an optimal value of N_c , above which they level off or even reverse. With increasing N_c , generally, one might expect a dramatic increase in the MC relaxation time due to an exponential increase in the state space. However, using the greedy MC algorithm, we only observe a gentle increase in both the number of MC sweeps and the CPU time while achieving excellent optimization results in terms of very low residual values of the cost function.

V. APPLICATIONS TO REAL DATA

A. Remotely sensed data

We investigate the application of the spin-based models on real-world data. We compare the results with the KNN, FKNN, and support vector machine (SVM) methods. We consider remotely sensed rainfall data on a moderately sized 50×50 grid. The study domain covers an area of Indonesia extending from 2.5S to 10N in latitude and from 110E to 122.5E in longitude, covered with a resolution of $0.25^{\circ} \times 0.25^{\circ}$. A map of the precipitation distribution is shown in Fig. 7(a). The data represent daily accumulated rainfall values recorded during January 2007 [44]. The values are in millimeters and some summary statistics are as follows: $z_{\min} = 7.1$, $z_{\max} = 832.6$, $\overline{z} = 295.0$, $z_{0.50} = 293.7$, and $\sigma_z = 151.4$. The skewness coefficient is 0.27 and the kurtosis coefficient is 2.93. As evidenced in the histogram shown in Fig. 7(b), the precipitation PDF is non-Gaussian, possibly bimodal.

The KNN and FKNN results are obtained with an optimally tuned parameter $k_{\rm opt}$ for each realization. In the case of SVM, two hyperparameters, C and σ_k , need to be tuned for each of the N_c binary classifiers. For different values of C, we found the bandwidth value $\sigma_{k_{\rm opt}}$ that minimizes globally the "testing errors" for each of the N_c classifiers (the true values at all the prediction points were used for testing). We selected a $C_{\rm opt}$ that minimizes the testing error. The obtained

TABLE VIII. Average misclassification rates for the N_c -level INNC, PNNC, and CNNC models. The averages are calculated over ten randomly thinned [by p (%)] samples of the rainfall data. These are compared with the best (lowest values) obtained using the KNN and SVM techniques. The following misclassification rates are used: $\langle F_I^* \rangle$ for the INNC, $\langle F_P^* \rangle$ for the PNNC, $\langle F_C^* \rangle$ for the CNNC, $\langle F_{\text{KNN}}^* \rangle$ for the KNN, and $\langle F_{\text{SVM}}^* \rangle$ for the SVM.

No. of classes	N_c	=2	N_c	=4	N_c	=8	N_c :	=16
p (%)	33	66	33	66	33	66	33	66
$\langle F_I^* \rangle$	6.81	8.61	13.41	16.94	24.51	30.46	44.95	53.05
$\langle F_P^* \rangle$	6.51	8.43	14.18	17.53	27.87	33.01	51.59	56.25
$\langle F_C^* \rangle$	6.86	8.62	14.68	17.84	28.63	32.97	50.32	55.28
$\langle F_{\mathrm{KNN}}^* \rangle$	6.46	8.47	13.82	18.01	28.28	33.85	51.85	56.12
$\langle F_{ m FKNN}^* \rangle$	6.12	7.91	12.99	16.72	26.55	31.63	49.67	54.08
$\langle F_{\mathrm{SVM}}^* \rangle$	6.78	8.56	15.61	17.62	29.70	33.61	51.27	55.49

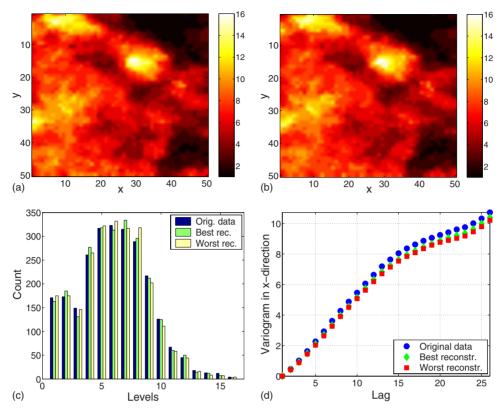


FIG. 8. (Color online) 16-class classification results of the rainfall data with p=33% missing values obtained by the INNC model. Using 100 reconstructions derived from different sampling configurations, we show the class maps of the best (a) and worst (b) case reconstructions, the class histograms of the original data as well as the best and worst reconstructions (c), and the empirical class variogram in the horizontal direction (d) for the original data as well as the best and worst reconstructions.

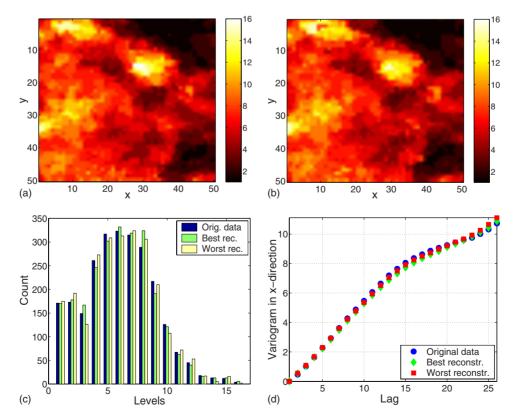


FIG. 9. (Color online) 16-class classification results of the rainfall data obtained by the PNNC model. The plot captions correspond to those in Fig. 8.

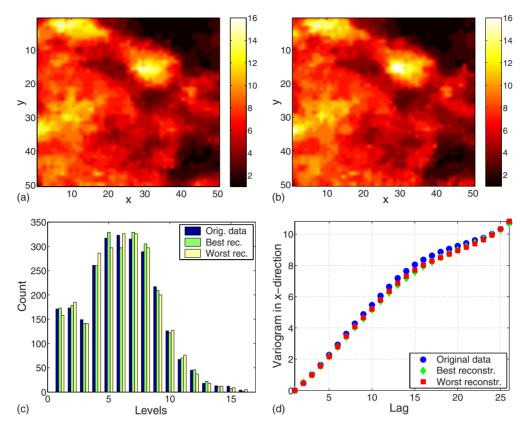


FIG. 10. (Color online) 16-class classification results of the rainfall data obtained by the CNNC model. The plot captions correspond to those in Fig. 8.

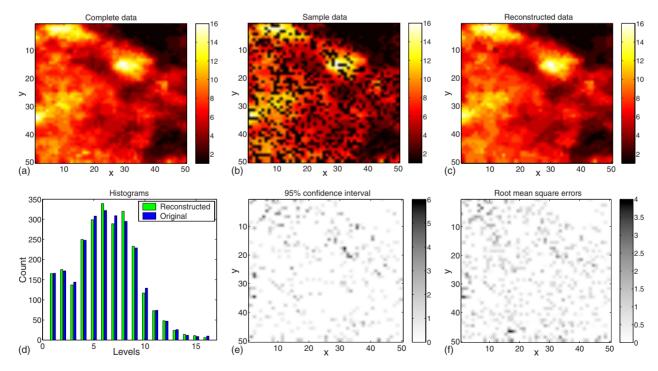


FIG. 11. (Color online) Classification results obtained from the 16-class CNNC model. The missing values sample is obtained from the rainfall data by randomly removing p=0.33 of the points. The 100 reconstructions are generated starting from 100 initial states obtained by the majority rule with adaptable stencil size. Plots (a)–(c) show class maps for the complete realization, the sample with the missing data, and the reconstructed image, respectively. The latter is based on the medians of the class values obtained from the 100 reconstructions. Plot (d) compares the class histograms of the original and reconstructed data, plot (e) shows the width of the 95% confidence intervals for the class predictions, and plot (f) represents the class root-mean-square error at the prediction points.



FIG. 12. Degraded image of Lena obtained by random removal of 33% of the pixels.

misclassification rates were relatively high in this class-adaptive approach, and further fine tuning of the hyperparameters did not bring noticeable improvements. Better results were obtained by using mean values of $\sigma_{\rm k_{opt}}$ and $C_{\rm opt}$ for all N_c classifiers. Using the same values of the hyperparameters for all the N_c classifiers has been shown to give satisfactory results [45]. Furthermore, we used the same hyperparameters for the ten realizations since they were all derived from a uniform thinning of the same rainfall data set.

In Table VIII we compare average misclassification rates for the spin-based models, as well as the KNN, FKNN, and SVM classifiers based on ten different sampling realizations. The sample sets are derived from random thinning of the rainfall data by p (%). Comparing the performance of different classifiers, the FKNN and the INNC models give the best results overall. FKKN performs better for small number of classes $(N_c=2,4)$, while INNC is better at larger values of N_c . These results agree with the synthetic data studies, where the INNC model also showed superior performance, especially for the data with slower variation and moderate degree of thinning. In contrast to the synthetic studies, the CNNC model did not perform as satisfactorily. This could be attributed to the presence of noise (ubiquitous in real data), as pointed out in Sec. IV. In addition, the SVM classifier did not perform as well as in some other comparative studies [29,30]. This may be due to the simplifications we adopted in



FIG. 13. Original (a) and reconstructed images of Lena using the INNC (b), PNNC (c), and CNNC (d) classification models.

the estimation of the hyperparameters. Besides the misclassification rates, we also checked the capacity of the respective classifiers to reproduce the histogram and the empirical directional variogram, $\hat{\gamma}_Z(r\mathbf{e})$ of the data [1]. The $\hat{\gamma}_Z(r\mathbf{e})$, also known as the two-point structure function, is a measure of the spatial continuity of the field Z in the direction $\mathbf{e} = \hat{\mathbf{x}}$. On a lattice of step α , the $\hat{\gamma}_Z(r\mathbf{e})$ is given by

$$\hat{\gamma}_{Z}(n\alpha\hat{\mathbf{x}}) = \frac{1}{2L(L-n)} \sum_{j=1}^{L-n} \sum_{i=1}^{L} \left[Z(\mathbf{s}_{i,j+n}) - Z(\mathbf{s}_{i,j}) \right]^{2},$$

$$n = 1, \dots, \frac{L}{2}, \tag{7}$$

where $\mathbf{s}_{i,j} = \alpha(i\hat{\mathbf{y}} + j\hat{\mathbf{x}})$. In Figs. 8–10, we show the reconstructed maps, histograms, and variograms in the direction of the *x* axis for the best (lowest F^*) and worst (highest F^*)

TABLE IX. Prediction errors and average optimization statistics for reconstructions of the Lena image by means of the INNC, PNNC, and CNNC models based on ten realizations corresponding to different sampling configurations. MAAE: mean average absolute error. MARE: mean average relative error. MAARE: mean average absolute relative error. MRASE: mean root average square error. MR: mean correlation coefficient. First, averages are evaluated over the prediction points for each realization, and then means are calculated with respect to an ensemble of ten realizations. The number in square brackets indicates the power of 10.

	MAAE	MARE (%)	MAARE (%)	MRASE	MR (%)	$\langle N_{ m MC} angle$	$\langle T_{\mathrm{CPU}} \rangle$ (s)	$\langle U^* angle$
INNC	5.04	2.63	4.97	9.29	98.35	106.9	351.1	3[-4]
PNNC	6.20	0.21	6.48	11.67	97.08	571.4	275.7	0.14
CNNC	5.04	-3[-6]	5.33	8.24	98.52	8.3	173.7	2[-10]

reconstructed realizations for N_c =16 and p=0.33 based on 100 realizations. The variogram along the y axis (not shown) has similar behavior. In all the cases, the statistics are recovered satisfactorily, and there are no significant differences between the respective models. Even though the PNNC model gave the highest misclassification rate, it reproduces the histograms and the variogram reasonably well.

The above simulation results are based on one run for each sample set. As we have shown in the case of synthetic data, multiple runs can improve the results and allow estimation of uncertainty. In the case of the rainfall data, considering one realization (sample set) generated with $p\!=\!0.33$ and performing 100 simulation runs using, for example, the 16-level CNNC model gave the misclassification rate F^* =46.9% requiring $T_{\text{CPU}}\!=\!5.6$ s of total CPU time. The multiple-run-reconstruction measures, as those shown in Fig. 6 for the synthetic data, are displayed in Fig. 11.

B. Digital image data

We consider the standard 256-valued gray-scale test image of Lena on a 512×512 grid. We randomly remove p = 33% of the pixels and subsequently reconstruct the image using the spin-based models. The degraded image is shown in Fig. 12.

The range of the image pixel values is equal to $\max(Z) - \min(Z) = 220$, and thus we use $N_c = 220$ classes of uniform width (equal to 1). The sequence in Fig. 13 shows the original image [Fig. 13(a)] along with the reconstructed images obtained by the INNC [Fig. 13(b)], PNNC [Fig. 13(c)], and CNNC [Fig. 13(d)] models. Visually, all three reconstructions appear quite similar to the original image. The INNC model misclassifies some pixels along the edges (e.g., along the shoulder). The numerical comparisons of univariate validation measures shown in Table IX are in favor of the CNNC model. The CNNC model is also more efficient computationally (requiring fewer MC sweeps and less CPU time), and it also approximates more accurately the sample energy. The worst performance in terms of the validation measures is shown by the PNNC model.

VI. SUMMARY AND CONCLUSIONS

We present spatial classification methods for missing data on partially sampled Cartesian grids based on nonparametric spin models from statistical physics (e.g., Ising, Potts, clock, and XY models). The methods are based on the idea of matching the normalized correlation energy of the sampled spins with that of the entire grid. The matching is performed using greedy Monte Carlo simulation conditioned by the sample values. The nonparametric spin-based classifiers presented here embody isotropic nearest-neighbor correlations. Hence, they are expected to perform well if the data exhibit a degree of spatial continuity and structure dominated by local features. Many geophysical and remotely sensed data as well as digital images share these features. The models presented here are not in principle suitable for capturing long-range correlations, such as those characterizing the transport properties of geological media. Nonetheless, it should be pointed out that the spatial features of the "reconstructed" field are not determined exclusively by the properties of the classification model but also by the features (e.g., anisotropic correlations) present in the sample.

The relative performance of the spin models in the case studies investigated varied, depending on the type of data, the sampling density, and the number of classes considered. Overall, the INNC model, which is based on a sequential classification algorithm, gave the most accurate classification rates in most of the cases studied. For all the simulated data, the PNNC model gave the highest misclassification rates. We believe that this is mainly due to the higher spatial degeneracy of the PNNC model. For noise-free data with shortrange differentiable variations, the CNNC model that incorporates cross-class correlations performed best, especially for the higher class numbers. As the number of classes increases, the CNNC model tends to the continuous XY model, and the classification emulates spatial interpolation. Up to a threshold, increasing the number of classes gradually lowers spatial prediction errors at the cost of a moderate increase in computing time. The classification performance of the spinbased methods can be further improved by executing multiple runs starting from different initial states. This strategy also permits an estimate of the classification uncertainty. The classification performance of the spin-based models was compared with the k-nearest-neighbor (KNN) and the fuzzy k-nearest-neighbor (FKNN) algorithms and (for the rainfall data) with the SVM classifier. For the synthetic data the INNC and CNNC models gave uniformly lower misclassification rates than the KNN and FKNN algorithms. For classification of real data into a small number of classes, the FKNN algorithm with optimized k was the most accurate of the classifiers tested.

All the spin-based models are computationally efficient. For the PNNC, CNNC, and XYNNC models, the mean CPU time ranged from 0.03 s (PNNC model, L=50, p=0.33, and N_c =8) to 5 s (CNNC model, L=200, p=0.66, and N_c =16). For the INNC model, the CPU time is generally higher due to the cost of determining the initial state at each level. In contrast, the time needed for the Monte Carlo relaxation is very short. Therefore, the resulting INNC CPU time varies almost linearly with the number of classes, and in our study it ranged from 0.08 to almost 14 s. We do not report CPU times for the KNN, FKNN, and SVM computations since in order to optimize the classification accuracy significant computational time was devoted to fine tuning the hyperparameters.

An advantage of the spin-based models with respect to the other classifiers tested is the lack of hyperparameters that need tuning by the user. Hence the classification procedure can be automated, and it provides competitive accuracy as well as computational efficiency. Compared to linear spatial interpolation algorithms (e.g., kriging), the spin-based classification methods present the advantages of computational efficiency and ability to handle non-Gaussian probability distributions at the expense of introducing discrete intervals in place of continuous values. A comparative study of the two approaches in the future could help us to quantify their relative merits.

Currently, the spin-based models are formulated on a regular grid. Hence, potential areas of application involve the compression of large images and the reconstruction of missing values (e.g., image pixels). Note that in light of the comments in Sec. III, the energy matching principle is not suitable for the refinement (resampling) of a regular grid, e.g., by doubling the spatial resolution. Extension to irregular sampling patterns is possible by defining a distance-dependent interaction strength $J_{i,j} = J_0 K(\vec{r}_{ij})$, where J_0 is arbitrary and $K(\vec{r}_{ij})$ is a normalized function of \vec{r}_{ij} , over a specified interaction neighborhood. Other potential extensions include extended-range interactions and/or "multipoint spin" correlations in the respective Hamiltonians. This could also help us to eliminate the spatial degeneracy evident in the present

models and provide more flexible measures of spatial dependence at the expense of concomitant increases in computational time and parametrization.

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